

Radiative Corrections for Lepton Scattering

Andrei Afanasev

Hampton University and Jefferson Lab

Presentation for Fermilab Neutrino Community, Feb.22, 2008

Collaborators: I. Akushevich, N. Merenkov, A. Ilyichev, *K. Joo*, *V. Burkert*,
S. Brodsky, C. Carlson, M. Vanderhaeghen, *G. Gilfoyle*



Main problem

Accelerated charge radiates

- ...While radiative corrections were the largest corrections to the data, and involved a considerable amount of computation, they were understood to a confidence level of 5% to 10% and did not significantly increase the total error in the measurements.

Henry W. Kendall

Nobel Lecture, December 8, 1990

- **Uncertainties in QED radiative corrections limit interpretability of precision experiments on electron-hadron scattering**



Plan of talk

Radiative corrections for electron scattering

- . Model-independent and model-dependent; soft and hard photons
- . Refined bremsstrahlung calculations
- . Two-photon exchange effects in the process $e^+p \rightarrow e^+p$

Rad. corrections for electroweak processes



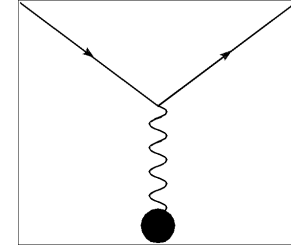
Example:

Measurements of Elastic Nucleon Form Factors

- Based on one-photon exchange approximation

$$M_{fi} = M_{fi}^{1\gamma}$$

$$M_{fi}^{1\gamma} = e^2 \bar{u}_e \gamma_\mu u_e \bar{u}_p (F_1(t) \gamma_\mu - \frac{\sigma_{\mu\nu} q_\nu}{2m} F_2(t)) u_p$$



- Two techniques to measure

$$\sigma = \sigma_0 (G_M^2 \tau + \varepsilon \cdot G_E^2) \quad : \text{Rosenbluth technique}$$

$$\frac{P_x}{P_z} = -\frac{A_x}{A_z} = -\frac{G_E \sqrt{\tau} \sqrt{2\varepsilon(1-\varepsilon)}}{G_M \tau \sqrt{1-\varepsilon^2}} \quad : \text{Polarization technique}$$

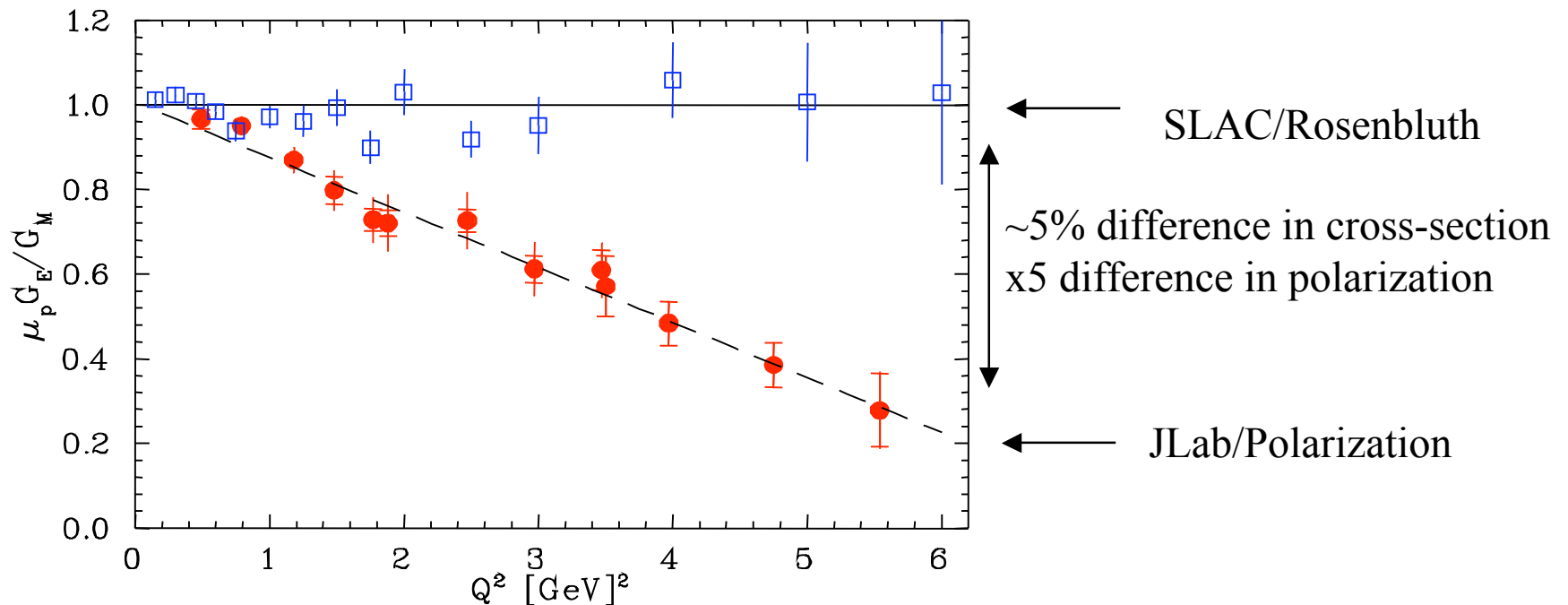
$$G_E = F_1 - \tau F_2, \quad G_M = F_1 + F_2$$

$$(P_y = 0)$$

Latter due to: Akhiezer, Rekalov; Arnold, Carlson, Gross



Do the techniques agree?

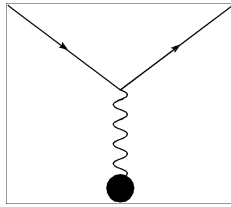


- Both early SLAC and Recent JLab experiments on (super)Rosenbluth separations followed $G_E/G_M \sim \text{const}$
- JLab measurements using polarization transfer technique give different results (Jones'00, Gayou'02)

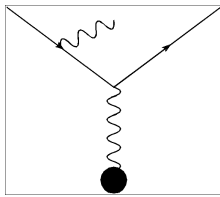
Radiative corrections, in particular, a short-range part of 2-photon exchange is a likely origin of the discrepancy



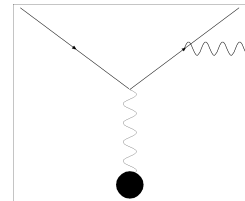
Basics of QED radiative corrections



(First) Born approximation

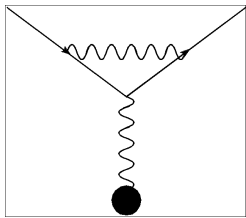


Initial-state radiation



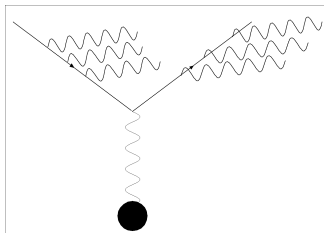
Final-state radiation

Cross section $\sim d\Omega/\Omega \Rightarrow$ integral diverges logarithmically: **IR catastrophe**



Vertex correction \Rightarrow cancels divergent terms; Schwinger (1949)

$$\sigma_{\text{exp}} = (1 + \delta)\sigma_{\text{Born}}, \quad \delta = \frac{-2\alpha}{\pi} \left\{ \left(\ln \frac{E}{\Delta E} - \frac{13}{12} \right) \left(\ln \frac{Q^2}{m_e^2} - 1 \right) + \frac{17}{36} + \frac{1}{2} f(\theta) \right\}$$

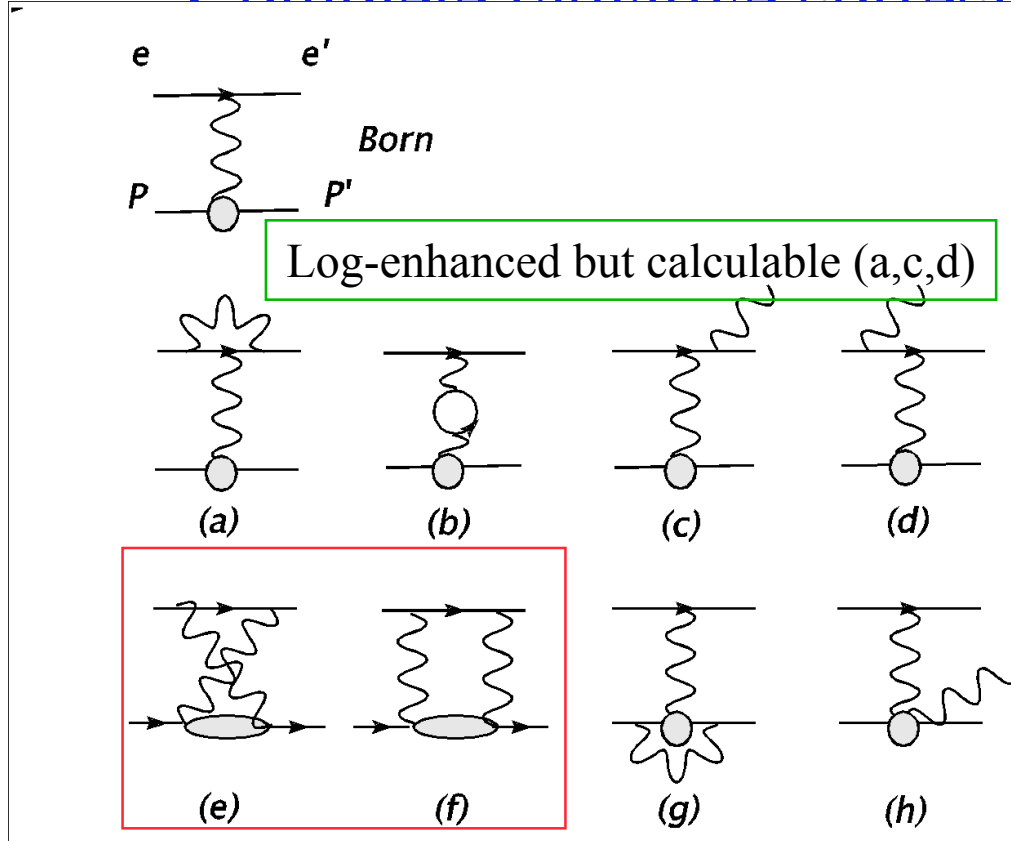


Multiple soft-photon emission: solved by exponentiation,
Yennie-Frautschi-Suura (YFS), 1961

$$(1 + \delta) \rightarrow e^\delta$$



Complete radiative correction in $O(\alpha_{em})$



Radiative Corrections:

- Electron vertex correction (a)
- Vacuum polarization (b)
- Electron bremsstrahlung (c,d)
- Two-photon exchange (e,f)
- Proton vertex and VCS (g,h)
- Corrections (e-h) depend on the nucleon structure
- Meister&Yennie; Mo&Tsai
- Further work by Bardin&Shumeiko; Maximon&Tjon; AA, Akushevich, Merenkov;
- Guichon&Vanderhaeghen'03:
Can (e-f) account for the Rosenbluth vs. polarization experimental discrepancy? Look for ~3% ...

Main issue: Corrections dependent on nucleon structure

Model calculations:

- Blunden, Melnitchouk, Tjon, Phys.Rev.Lett.**91**:142304,2003
- Chen, AA, Brodsky, Carlson, Vanderhaeghen, Phys.Rev.Lett.**93**:122301,2004



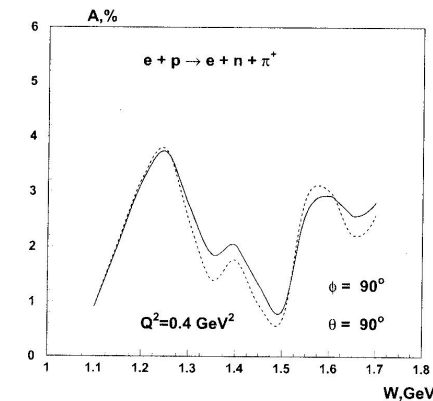
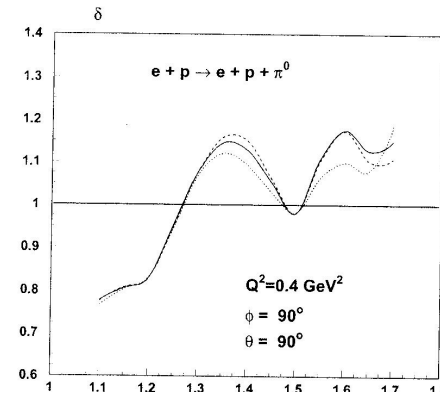
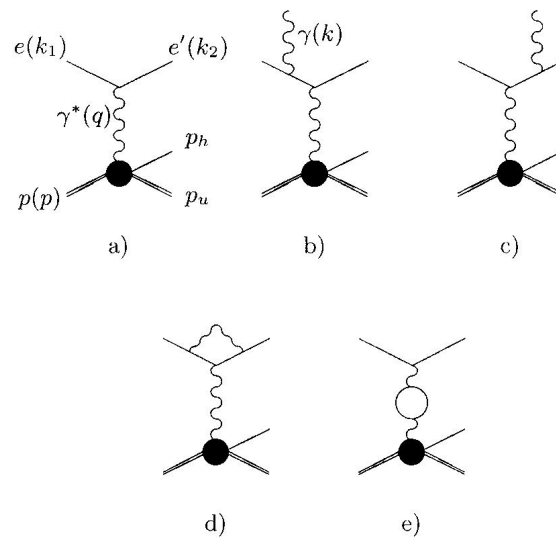
Basic Approaches to QED Corrections

- L.W. Mo, Y.S. Tsai, Rev. Mod. Phys. 41, 205 (1969); Y.S. Tsai, Preprint SLAC-PUB-848 (1971).
 - Considered both elastic and inelastic inclusive cases. No polarization.
- D.Yu. Bardin, N.M. Shumeiko, Nucl. Phys. B127, 242 (1977).
 - Covariant approach to the IR problem. Later extended to inclusive, semi-exclusive and exclusive reactions with polarization.
- E.A. Kuraev, V.S. Fadin, Yad.Fiz. 41, 7333 (1985); E.A. Kuraev, N.P.Merenkov, V.S. Fadin, Yad. Fiz. 47, 1593 (1988).
 - Developed a method of electron structure functions based on Drell-Yan representation; currently widely used at e^+e^- colliders.



RC for Electroduction of Pions

- **AA, Akushevich, Burkert, Joo, Phys.Rev.D66, 074004 (2002)**
 - Conventional RC, precise treatment of phase space, no peaking approximation, no dependence on hard/soft photon separation
 - Can be used for any exclusive electroproduction of 2 hadrons, e.g., $d(e,e'p)n$ (EXCLURAD code)



See <http://www.jlab.org/RC> for other codes
 Used in data analysis at JLab
 (and MIT, HERMES, MAMI,...)



Bethe-Heitler corrections to polarization transfer and cross sections

AA, Akushevich, Merenkov Phys.Rev.D64:113009,2001;
AA, Akushevich, Ilychev, Merenkov, PL B514, 269 (2001)

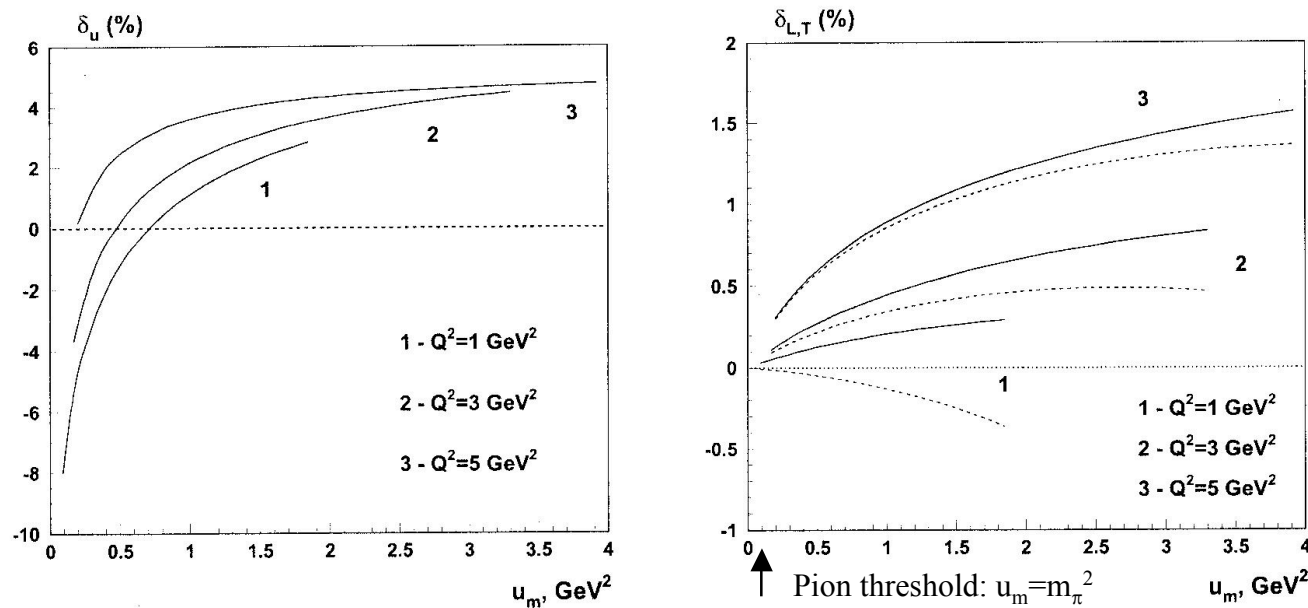


Figure 2: Radiative corrections to the unpolarized cross section (left plot) and polarization asymmetries (right plot) defined in (41). Solid and dashed lines corresponds to longitudinal and transverse cases. $S=8$ GeV².

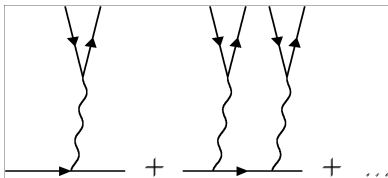
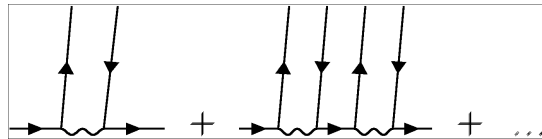
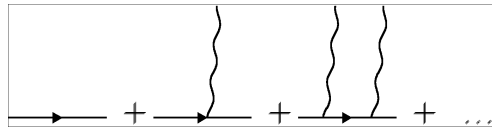
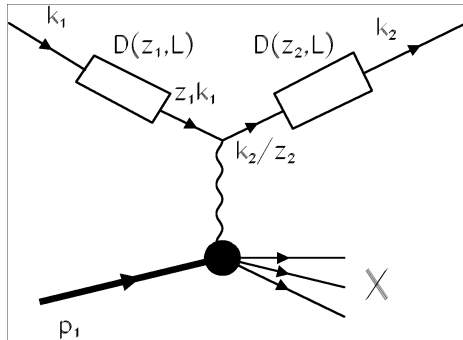
In kinematics of elastic ep-scattering measurements, cross sections are more sensitive to RC



Electron Structure Functions

(Kuraev,Fadin,Merenkov)

- For polarized $ep \rightarrow e'X$ scattering, AA et al, JETP 98, 403 (2004); elastic ep : AA et al. PRD 64, 113009 (2001).

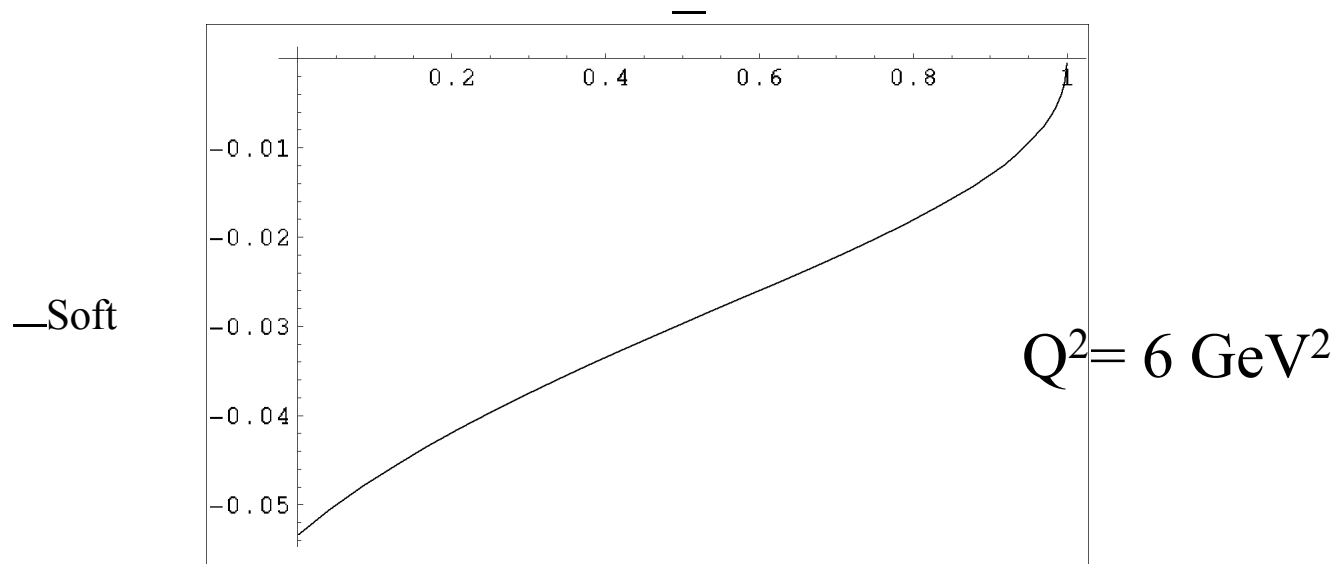


- Resummation technique for collinear photons (=peaking approx.)
- Difference $< 0.5\%$ from previous calculation including hard brem

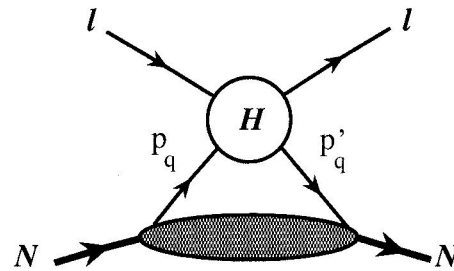


Separating *soft* 2-photon exchange

- Tsai; Maximon & Tjon (k_0)
- Grammer & Yennie prescription PRD 8, 4332 (1973) (also applied in QCD calculations)
- Shown is the resulting (soft) QED correction to [cross section](#)
- **Already included in experimental data analysis**
- **NB:** Corresponding effect to polarization transfer and/or asymmetry is zero

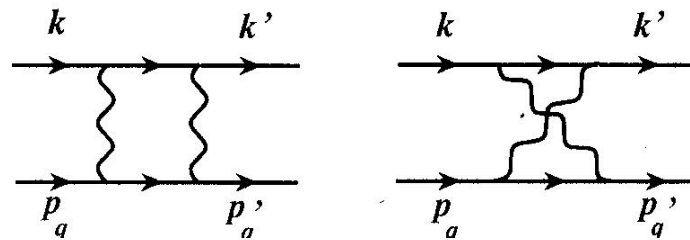


Calculations using Generalized Parton Distributions



Model schematics:

- Hard eq-interaction
- GPDs describe quark emission/absorption
- Soft/hard separation
 - Use Grammer-Yennie prescription



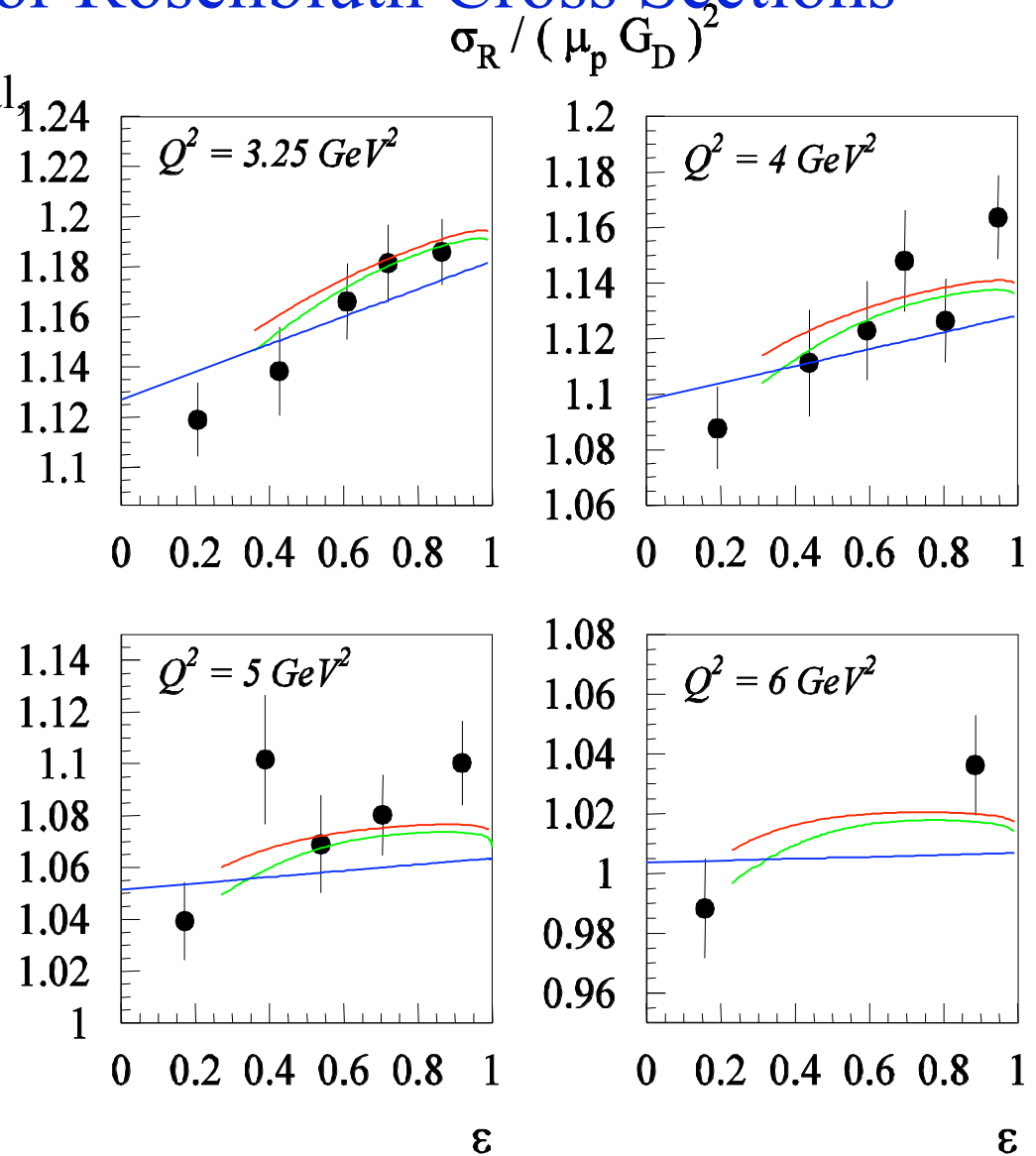
Hard interaction with a quark

AA, Brodsky, Carlson, Chen, Vanderhaeghen,
Phys.Rev.Lett.**93**:122301,2004; Phys.Rev.D**72**:013008,2005



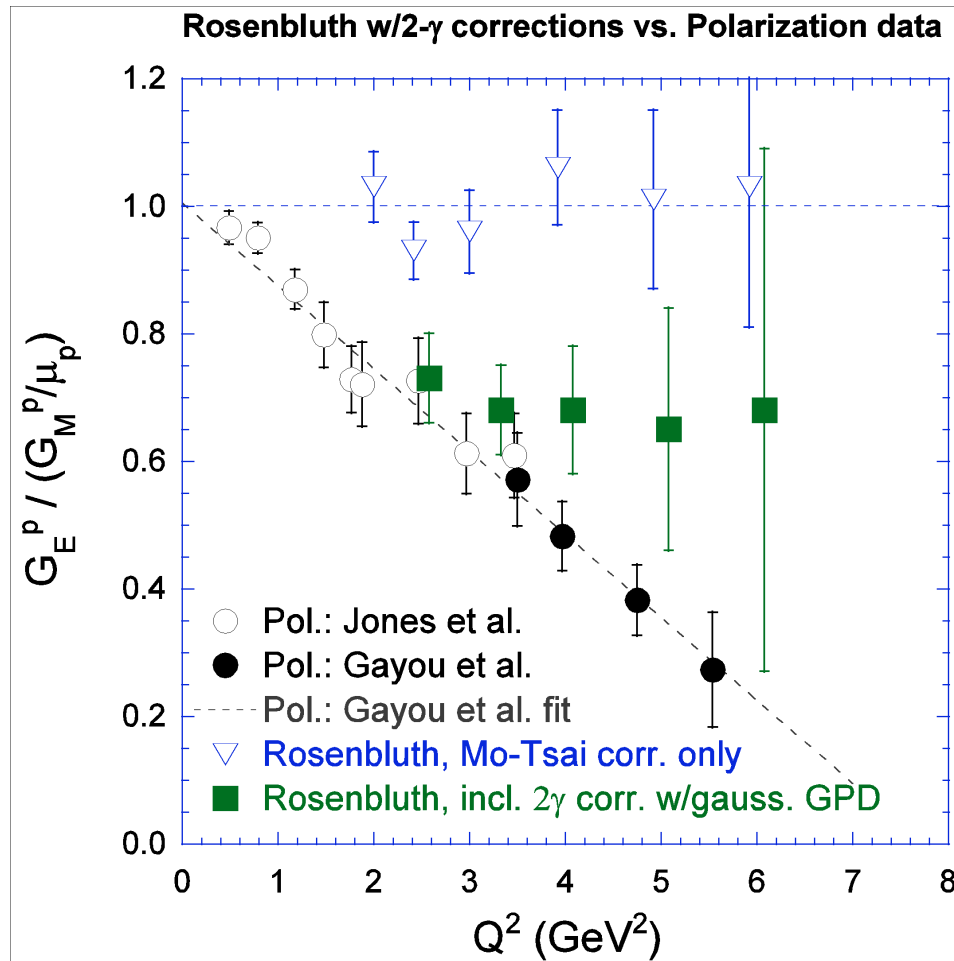
Two-Photon Effect for Rosenbluth Cross Sections

- Data shown are from Andivahis et al. PRD 50, 5491 (1994)
- Included GPD calculation of two-photon-exchange effect
- Qualitative agreement with data:
- Discrepancy likely reconciled



Updated Ge/Gm plot

AA, Brodsky, Carlson, Chen, Vanderhaeghen,
Phys.Rev.Lett.93:122301, 2004; Phys.Rev.D72:013008, 2005



Full Calculation of Bethe-Heitler Contribution

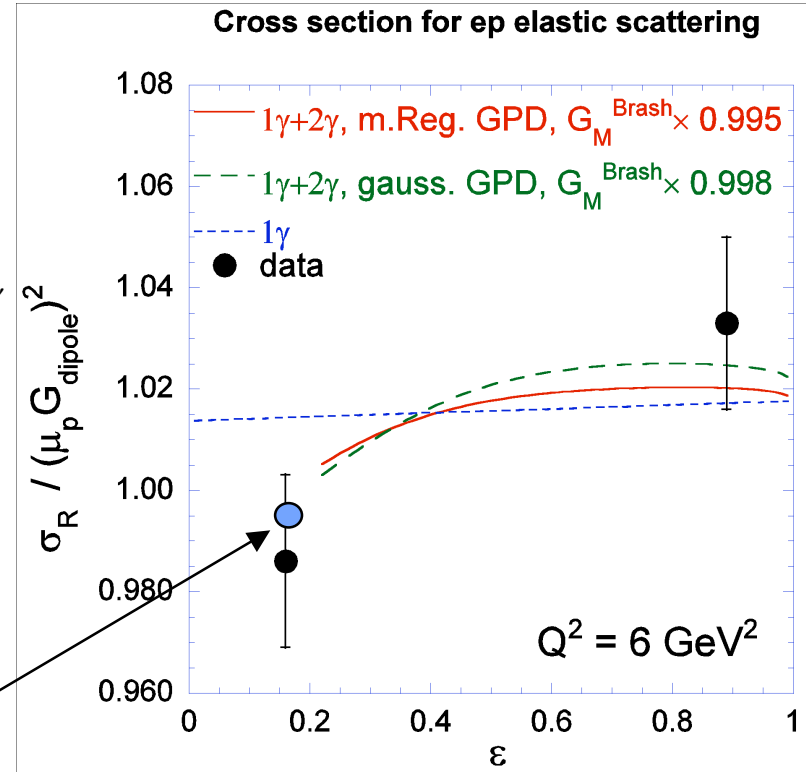
*Additional work by AA et al., using MASCARAD (Phys.Rev.D64:113009,2001)
Full calculation including soft and hard bremsstrahlung*

Radiative leptonic tensor in full form
AA et al, **PLB 514, 269 (2001)**

$$L_{\mu\nu}^r = -\frac{1}{2} \text{Tr}(\hat{k}_2 + m) \Gamma_{\mu\alpha} (1 + \gamma_5 \hat{\xi}_e) (\hat{k}_1 + m) \bar{\Gamma}_{\alpha\nu}$$

$$\Gamma_{\mu\alpha} = \left(\frac{k_{1\alpha}}{k \cdot k_1} - \frac{k_{2\alpha}}{k \cdot k_2} \right) \gamma_\mu - \frac{\gamma_\mu \hat{k} \gamma_\alpha}{2k \cdot k_1} - \frac{\gamma_\alpha \hat{k} \gamma_\mu}{2k \cdot k_2}$$

$$\Gamma_{\alpha\nu} = \left(\frac{k_{1\alpha}}{k \cdot k_1} - \frac{k_{2\alpha}}{k \cdot k_2} \right) \gamma_\nu - \frac{\gamma_\alpha \hat{k} \gamma_\nu}{2k \cdot k_1} - \frac{\gamma_\nu \hat{k} \gamma_\alpha}{2k \cdot k_2}$$

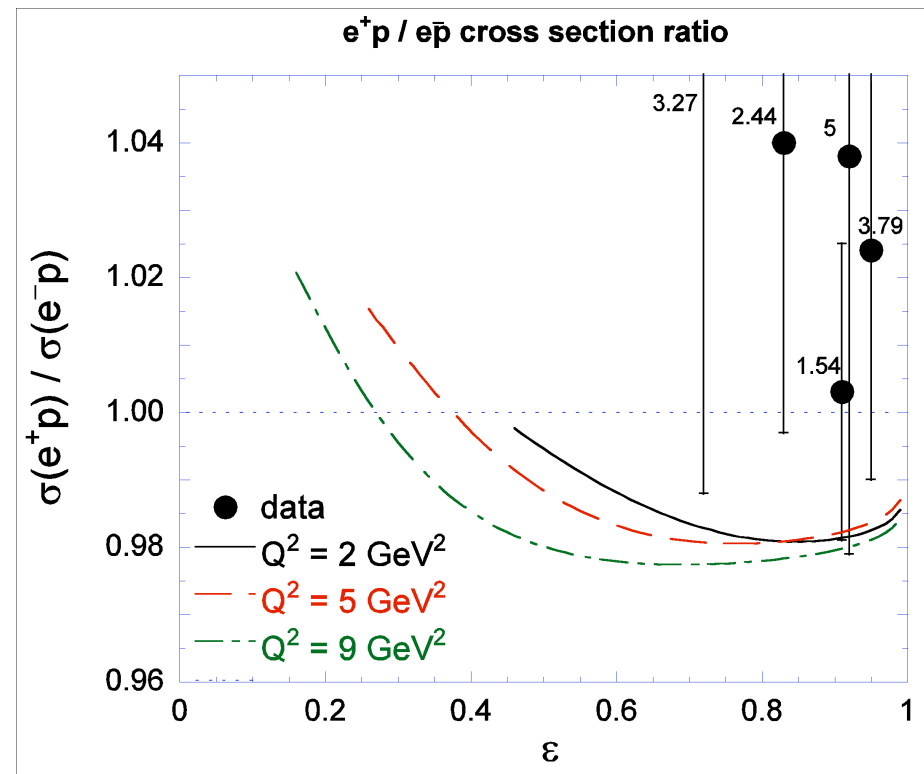


Additional effect of full soft+hard brem $\rightarrow +1.2\%$ correction to ϵ -slope
Resolves additional ~25% of Rosenbluth/polarization discrepancy!



Charge asymmetry

- Cross sections of electron-proton scattering and positron-proton scattering are equal in one-photon exchange approximation
 - Different for two- or more photon exchange



To be measured in JLab Experiment 04-116,
Spokepersons AA, W. Brooks, L.Weinstein, et al.

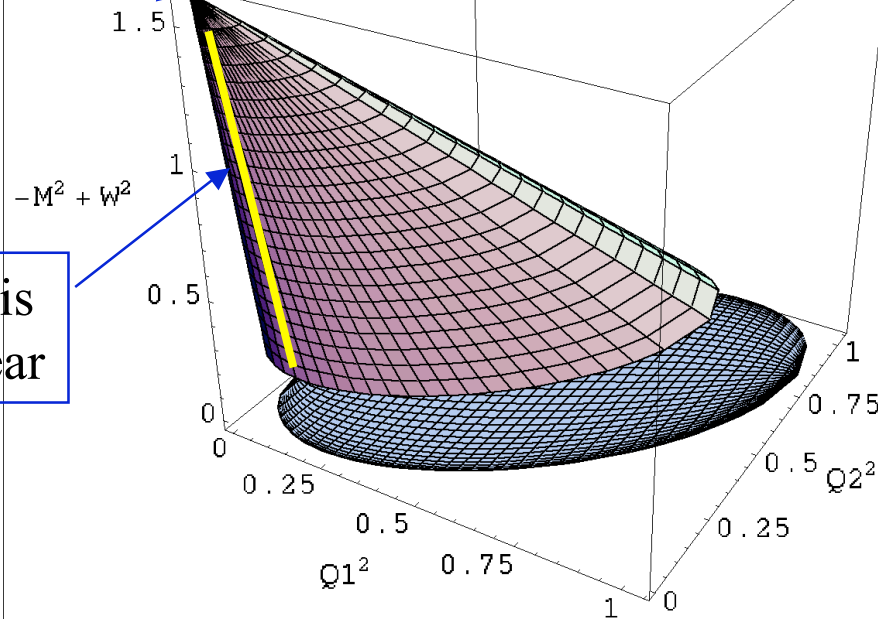


Phase Space Contributing to the absorptive part of 2_{γ} -exchange amplitude

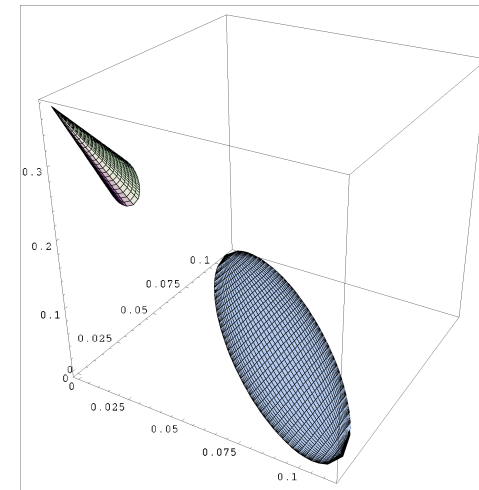
- 2-dimensional integration (Q_1^2, Q_2^2) for the elastic intermediate state
- 3-dimensional integration (Q_1^2, Q_2^2, W^2) for inelastic excitations

'Soft' intermediate electron;
Both photons are hard collinear

One photon is
Hard collinear



Examples: MAMI A4
E= 855 MeV
 $\theta_{cm}= 57$ deg;
SAMPLE, E=200 MeV



Other theoretical developments

- Blunden et al., Phys.Rev.C72:034612, 2005
Approximate proton Compton amplitude by Born terms
- Kondratyuk et al., nucl-th/0506026
Add intermediate Δ -excitation to the above
- Pascalutsa et al., hep-ph/0509055
GPD approach extended to $N \Delta$ transition
- Borisjuk, Kobushkin, Phys.Rev.C72:035207,2005

Future task: Resummation of inelastic excitations at lower Q^2



Two-photon exchange for electron-proton scattering

- Quark-level short-range contributions are substantial (3-4%) ; correspond to $J=0$ fixed pole (Brodsky-Close-Gunion, PRD 5, 1384 (1972)).
- Structure-dependent radiative corrections calculated using GPDs bring into agreement the results of polarization transfer and Rosenbluth techniques for G_{ep} measurements
- Experimental tests of two-photon exchange
 - Comparison between electron and positron elastic scattering (JLab E04-116)
 - Measurement of nonlinearity of Rosenbluth plot (JLab E05-017)
 - Search for deviation of angular dependence of polarization and/or asymmetries from Born behavior at fixed Q^2 (JLab E04-019)
 - Elastic single-spin asymmetry or induced polarization (JLab E05-015)
 - 2_+ additions for parity-violating measurements (HAPPEX, G0)

Through active theoretical support emerged a research program of
Testing precision of the electromagnetic probe
Double-virtual VCS studies with two space-like photons



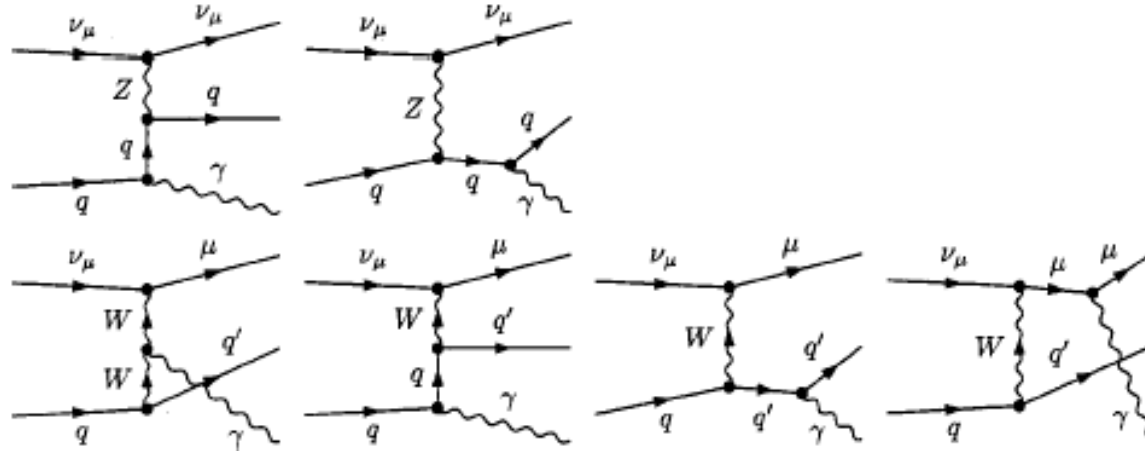
Radiative Corrections for Electro-Weak Processes

Semi-Leptonic processes involving nucleons

- Neutrino-nucleon scattering
 - Per cent level reached by NuTeV. Radiative corrections for DIS calculated at a partonic level (D. Bardin et al.)
- Neutron beta-decay: Important for V_{ud} measurements; axial-vector coupling g_A
Marciano, Sirlin, PRL 56, 22 (1986); Ando et al., Phys.Lett.B595:250-259,2004; Hardy, Towner, PRL94:092502,2005
 - Extended to $_N$ by Fukugita, Acta Phys.Polon.B35:1687-1732,2004; and $_D$: Phys.Rev.D72:071301,2005, Erratum-ibid.D74:039906,2006;
 - Kurylov, Phys.Rev. C65 (2002) 055501= \Rightarrow $\sim 4\%$ effect for $_{total}$
- Parity-violating DIS: Bardin, Fedorenko, Shumeiko, Sov.J.Nucl.Phys.32:403,1980; J.Phys.G7:1331,1981, up to 10% effect from rad.corrections
- Parity-violating elastic ep (strange quark effects, weak mixing angle)



Implications for Nutev



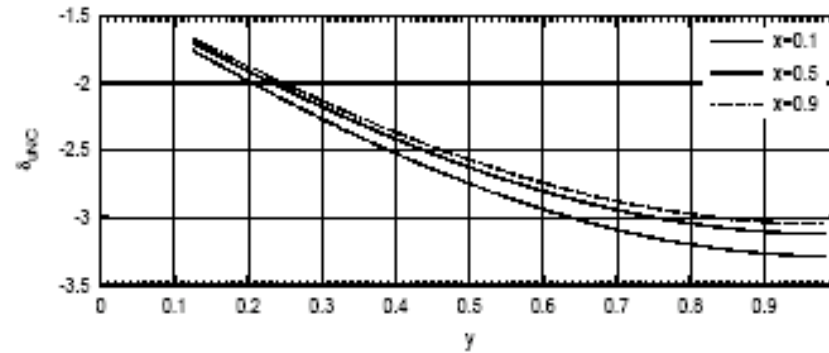
- Diener, Dittmaier, Hollik, Phys.Rev.D69:073005,2004:
 - **Rad.Corrections used by NuTeV likely underestimated,**
“...we compare results that differ in the input-parameter scheme, treatment of real photon radiation, and factorization scheme. The associated shifts in the theoretical prediction for the ratio of neutral- and charged-current cross sections can be larger than the experimental accuracy of the NuTeV result. ...”



Neutrino DIS

- Arbuzov, Bardin, Kalinovskaya, JHEP06(2005)078

$\delta_{uNC}(\%)$



$E_{\nu}=80\text{GeV}$

Figure 1: Relative effect of radiative corrections to $\nu - u$ NC scattering as a function of y for three fixed values of x .

$\delta_{dCC}(\%)$

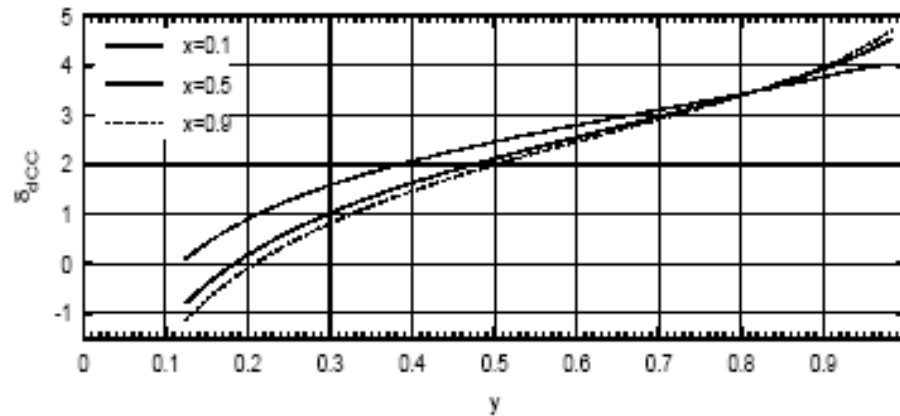


Figure 2: Relative effect of radiative corrections to $\nu - d$ CC scattering as a function of y for three fixed values of x .



Parity Violating elastic e-N scattering

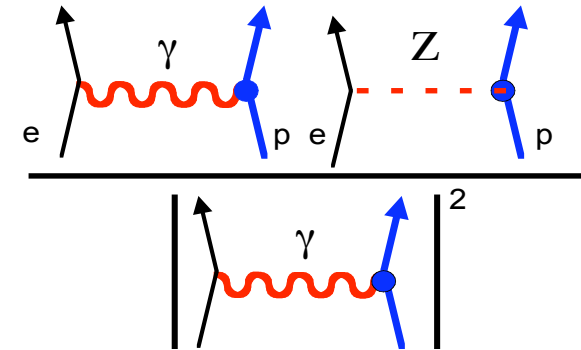
Longitudinally polarized electrons,
unpolarized target

$$A = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} = \left[\frac{-G_F Q^2}{4\pi\alpha\sqrt{2}} \right] \frac{A_E + A_M + A_A}{2\sigma_{unpol}}$$

$$A_E = \varepsilon(\theta) G_E^Z G_E^\gamma$$

$$A_M = \tau G_M^Z G_M^\gamma$$

$$A_A = -(1 - 4\sin^2 \theta_W) \varepsilon' G_A^e G_M^\gamma$$



$$\tau = Q^2/4M^2$$

$$\varepsilon = [1 + 2(1 + \tau)\tan^2(\theta/2)]^{-1}$$

$$\varepsilon' = [\tau(\tau + 1)(1 - \varepsilon^2)]^{1/2}$$

Neutral weak form factors contain explicit contributions from strange sea

$$G_{E,M}^Z(Q^2) = (1 - 4\sin^2 \theta_W)(1 + R_A^p)G_{E,M}^p - (1 + R_A^n)G_{E,M}^n - G_{E,M}^s$$

$$G_A^e(Q^2) = -G_A^Z + (\eta F_A^\gamma + R^e) + \Delta s$$

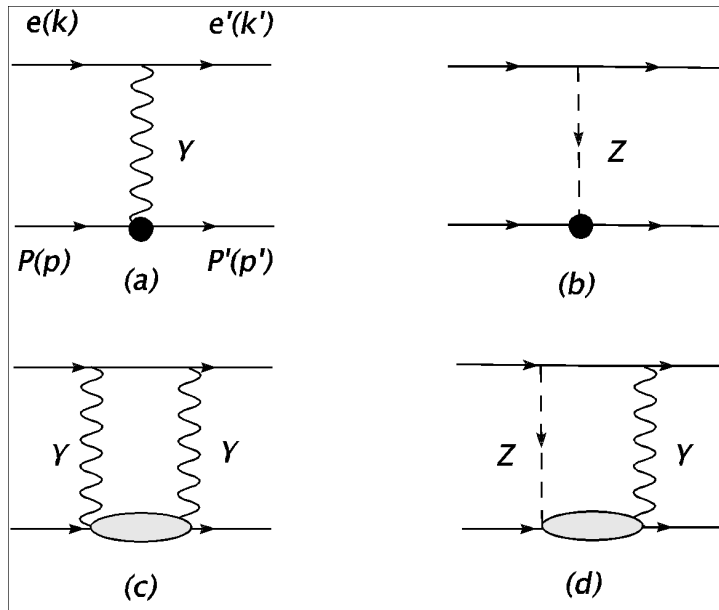
$$\eta = \frac{8\pi\alpha\sqrt{2}}{1 - 4\sin^2 \theta_W} = 3.45$$

$$G_A^Z(0) = 1.2673 \pm 0.0035 \text{ (from } \beta \text{ decay)}$$

Andrei Afanasev, Discussion with Fermilab Neutrino Community., 2/22/08



Born and Box diagrams for elastic ep-scattering

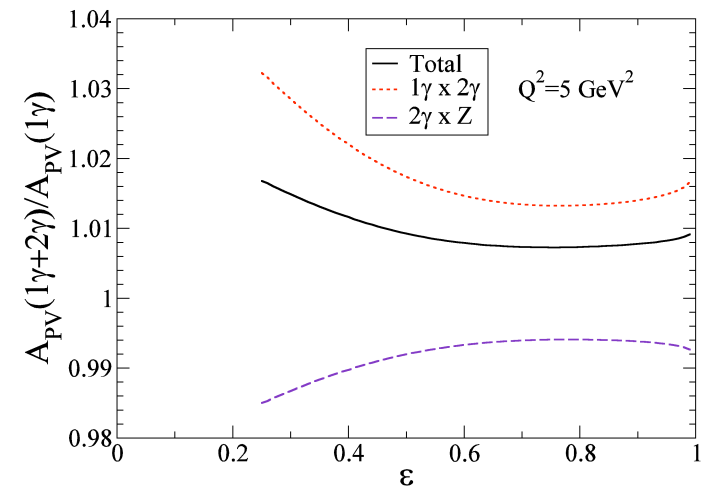
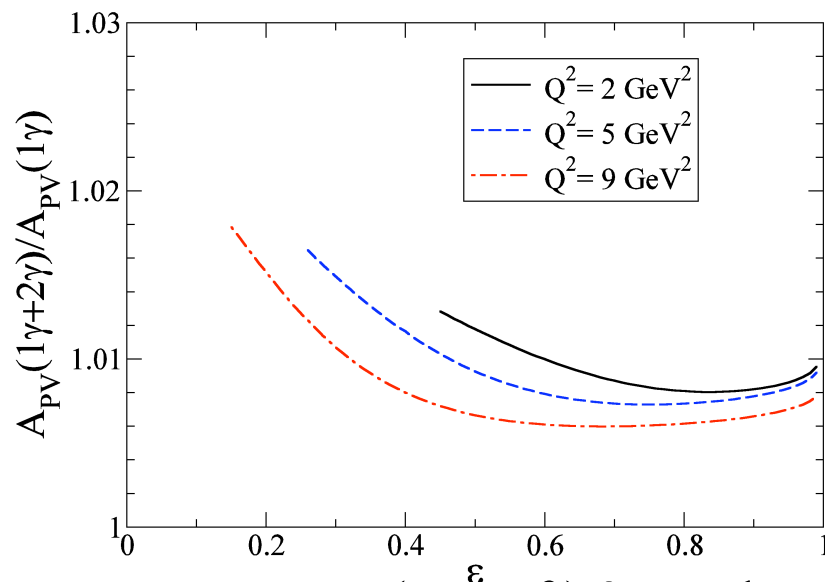


- (d) Computed by Marciano, Sirlin, Phys.Rev.D29:75,1984, Erratum-
ibid.D31:213,1985 for atomic PV (i.e., $Q^2 \rightarrow 0$)
- (c) Presumed small, e.g., M. Ramsey-Musolf, Phys.Rev. C60 (1999) 015501



GPD Calculation of 2_{γ} -Z-interference

- Can be used at higher Q^2 , but points at a problem of additional systematic corrections for parity-violating electron scattering. The effect evaluated in GPD formalism is the largest for backward angles:



AA & Carlson, *Phys. Rev. Lett.* **94**, 212301 (2005):

Measurements of strange-quark content of the nucleon are affected, s may shift by $\sim 10\%$

Important note: (nonsoft) 2_{γ} -exchange amplitude has no $1/Q^2$ singularity; 1_{γ} -exchange is $1/Q^2$ singular \Rightarrow At low Q^2 , 2_{γ} -corrections is suppressed as Q^2 P. Blunden used this formalism and evaluated correction of 0.16% for



2_-correction for ep-scattering via Z-exchange

- 2_-correction to parity-violating asymmetry **does not cancel**. May reach a few per cent for GeV momentum transfers
- Corrections are angular-dependent, not reducible to re-definition of coupling constants
 - Revision of γ -Z-box contribution and extension of model calculations to lower Q^2 is necessary
- Further developments:
 - Zhu, Kao, Yang, Phys.Rev.Lett.99:262001,2007: Found essential Q^2 -dependence of EW box contributions
 - Tjon, Melnitchouk, arXiv:0711.0143 [nucl-th]: Model calculation of EW box



RC for Minerva

- For CC cross sections, anticipate 1-5% electromagnetic effects
 - Bremsstrahlung calculations model-independent, but need to be matched with experimental cuts and acceptances
 - Electroweak box diagrams calculations depend on the used model of hadronic structure; can be constrained by existing (and forthcoming) info on 2_{γ} -exchange for elastic ep-scattering
- *Expertise at JLab available to implement Rad. Corrections for data analysis of Minerva*

